

(NB: presume existence)

Laplace's equation has the radial symmetric solution

$$\begin{cases} r^{2-n}, & n \geq 3 \\ \log r, & n = 2 \end{cases}$$

where  $r$  is the radial distance from a fixed point.

We thus fix a point  $y \in \Omega$  and introduce the normalised fundamental solution to the

Laplace equation:

$$P(x, y) = P(|x-y|) = \begin{cases} \frac{1}{n(2-n)\omega_n} |x-y|^{2-n}, & n \geq 3 \\ \frac{1}{2\pi} \log |x-y|, & n = 2 \end{cases}$$

where  $\omega_n = \frac{2\pi^{n/2}}{n\Gamma(\frac{n}{2})}$  = volume of the  $n$ -dim unit ball.

By a direct computation we have

$$P_i: P(x, y) = \frac{1}{n\omega_n} (x_i - y_i) |x-y|^{-n};$$

$$P_{ij}: P(x, y) = \frac{1}{n\omega_n} (|x-y|^{-2} \delta_{ij} - n(x_i - y_i)(x_j - y_j) |x-y|^{-n-2})$$

The singularity at  $x=y$  prevents us from using  $P$  in Green's second identity.

To overcome this we replace  $\Omega$  by  $\Omega \setminus \bar{B}_\rho$ , where  $B_\rho = B_\rho(y)$  for sufficiently small  $\rho > 0$ .

We can then conclude

$$\int_{\Omega \setminus B_\rho} \nabla u \cdot \nabla \phi \, dx = \int_{\partial \Omega} (\phi \frac{\partial u}{\partial \nu} - u \frac{\partial \phi}{\partial \nu}) \, d\sigma + \int_{\partial B_\rho} (\phi \frac{\partial u}{\partial \nu} - u \frac{\partial \phi}{\partial \nu}) \, d\sigma \quad (*)$$

Now

$$\int_{\partial B_\rho} \phi \frac{\partial u}{\partial \nu} \, d\sigma = P(\rho) \int_{\partial B_\rho} \frac{\partial u}{\partial \nu} \, d\sigma$$

$$\leq u \omega_n \rho^{n-1} P(\rho) \sup_{B_\rho} |\Delta u| \rightarrow 0 \text{ as } \rho \rightarrow 0,$$

and

$$\int_{\partial B_\rho} u \frac{\partial \phi}{\partial \nu} \, d\sigma = -P'(\rho) \int_{\partial B_\rho} u \, d\sigma \quad (\text{NB: } \nu \text{ is outer normal to } \Omega \setminus B_\rho)$$

$$= \frac{-1}{n\omega_n \rho^{n-1}} \int_{\partial B_\rho} u \rightarrow -u(y) \text{ as } \rho \rightarrow 0.$$

Thus by taking  $\delta \rightarrow 0$  in (†) we arrive at  
Green's representation formula :

$$u(y) = \int_{\partial\Omega} u \frac{\partial P}{\partial \nu}(x-y) - P(x-y) \frac{\partial u}{\partial \nu} d\sigma + \int_{\Omega} P(x-y) \Delta u dx, \quad y \in \Omega. \quad (\ddagger)$$

NB: for integrable  $f$ , the expression

$$\int_{\Omega} P(x-y) f(x) dx$$

is called the Newton potential with density  $f$ .

Remark. From (†) we can solve for  $u$  provided we know that

(i)  $\Delta u$  in  $\Omega$

(ii)  $u, \frac{\partial u}{\partial \nu}$  on  $\partial\Omega$ .

However the normal derivative  $\frac{\partial u}{\partial \nu}$  is unknown to us.

Remark. Suppose  $P(x, y) = P(x-y)$  solves  $\Delta P(x, y) = -\delta(x-y)$  on  $\mathbb{R}^n$

Then  $u(x) = \int P(x, y) f(y) dy$ , for given  $f$ , solves

$$\Delta_x u(x) = \int (\Delta_x P(x, y)) f(y) dy = \int \delta(x-y) f(y) dy = f(x)$$

on  $\mathbb{R}^n$ .

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To account for this, consider  $u \in C^1(\bar{\Omega}) \cap C^2(\Omega)$  satisfying

$$\Delta u = 0 \quad \text{in } \Omega.$$

Then by Green's second identity we obtain

$$-\int_{\partial\Omega} \left( u \frac{\partial h}{\partial \nu} - h \frac{\partial u}{\partial \nu} \right) d\sigma = \int_{\Omega} h \Delta u dx \quad (\S)$$

Now write  $G = P+h$  and add (†) with (‡) to obtain the more general version of the Green's representation formula:

$$u(y) = \int_{\partial\Omega} \left( u \frac{\partial G}{\partial \nu} - G \frac{\partial u}{\partial \nu} \right) d\sigma + \int_{\Omega} G \Delta u dx$$

Now if in addition  $G=0$  on  $\partial\Omega$  we obtain

$$u(y) = \int_{\partial\Omega} u \frac{\partial G}{\partial \nu} d\sigma + \int_{\Omega} G \Delta u dx$$

and the function  $G(x, y) = G$  is called the (Dirichlet) Green's function for the domain  $\Omega$ .

It gives a representation for  $C^1(\bar{\Omega}) \cap C^2(\Omega)$ -harmonic functions in terms of its boundary values.

(†) i.e.  $P+h=0$  on  $\partial\Omega$  so that  $h$  solves

$$\begin{cases} \Delta h = 0 & \text{in } \Omega \\ h = -P & \text{on } \partial\Omega. \end{cases}$$

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### § Physical meaning of Green's function

$$\begin{cases} \Delta_x G(x, y) = -\delta(x-y), & x \in \Omega \\ G = 0, & \text{on } \partial\Omega \end{cases}$$

$G(x, y)$  = Coulomb potential generated inside the grounded conducting surface  $\partial\Omega$  by the charge of  $+\frac{1}{4\pi}$  at the point  $y \in \Omega$ .

