

10TH SEPTEMBER, 2024

MODELLING MOLECULAR BEHAVIOUR IN A STRONG MAGNETIC FIELD BY EXPLOITING SPIN-SYMMETRY BREAKING

Emiel Vanden Berghe

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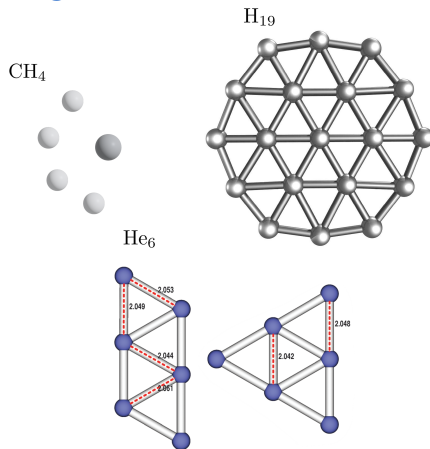
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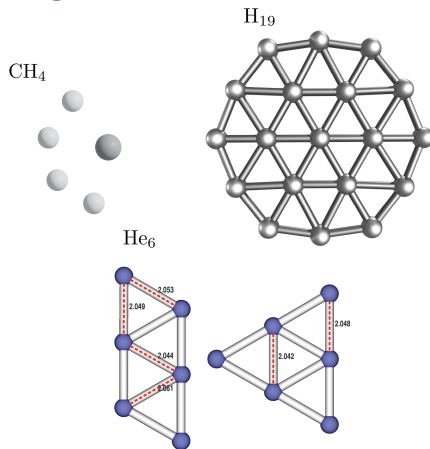
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- Spin-transitions to high-spin states associated with a **chemical regime-shift**:
 - Pauli exclusion **prohibits covalent bonding** between same spin electrons
 - Novel bonding mechanism
- Spin-transitions can **affect dissociation behaviour** even before the high-spin regime fully takes hold

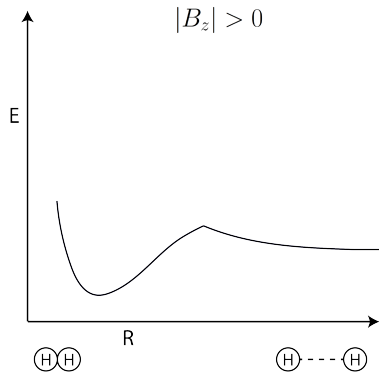
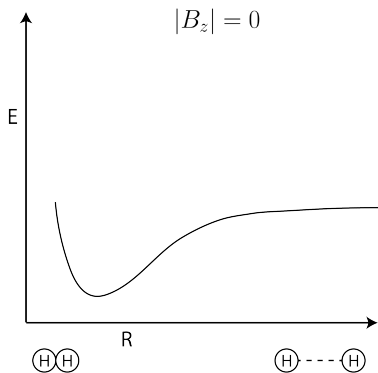


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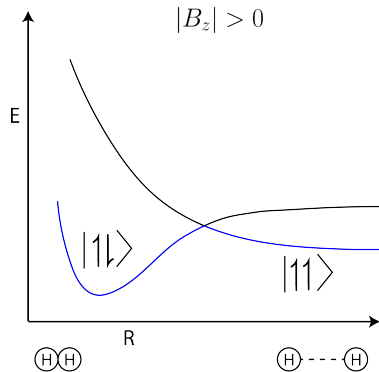
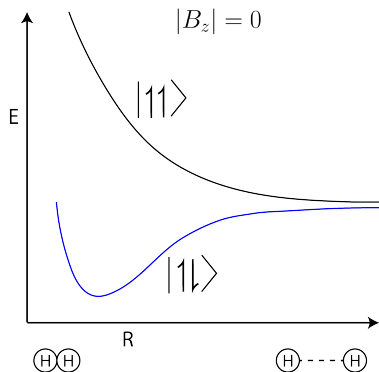
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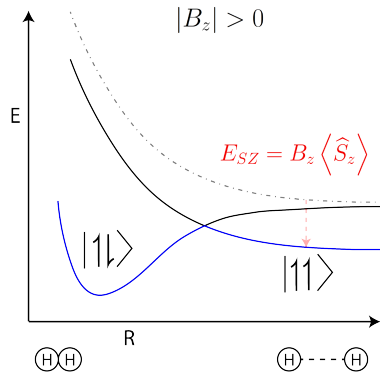
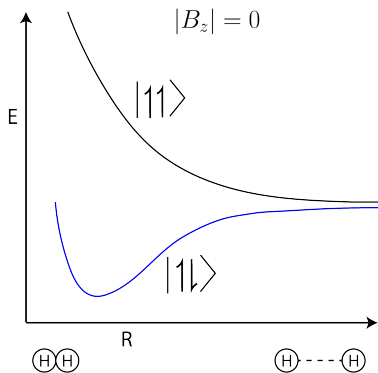
SPIN-TRANSITIONS THROUGHOUT DISSOCIATION: H_2 IN A MAGNETIC FIELD



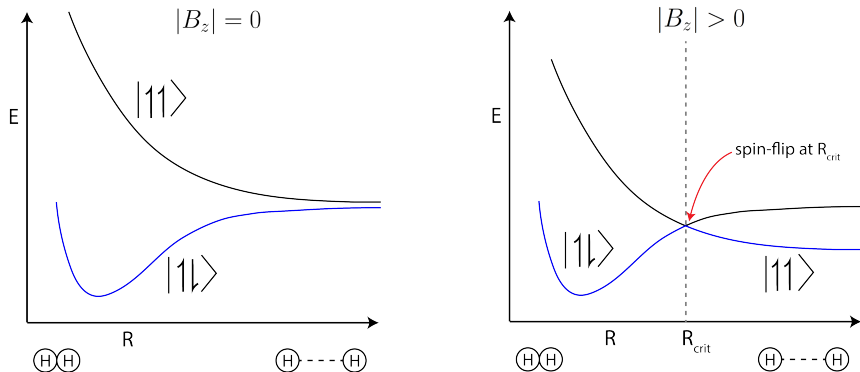
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Spin-transitions cause unexpected discontinuities in dissociation profiles

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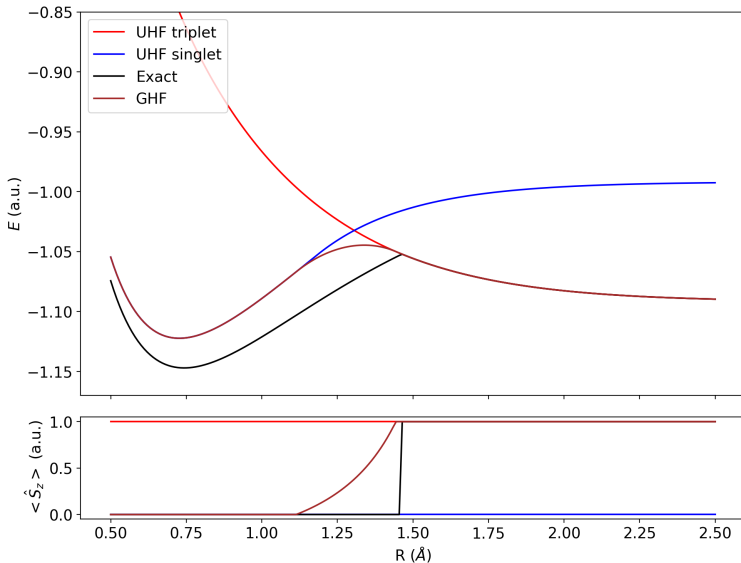
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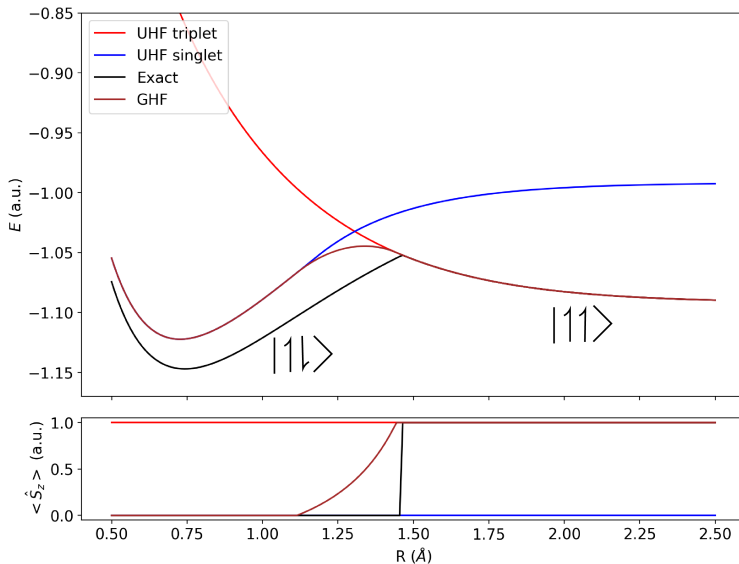
- Variational (non-perturbative):
 - Don't need field-free reference state
 - Model **non-linear effects** in electronic structure
 - *Minimal-coupling*

$$\hat{\mathcal{H}}(\mathbf{B}) = \hat{\mathcal{H}}_0 + \underbrace{\mathbf{B} \cdot \hat{\mathbf{S}}}_{\text{Spin-Zeeman}} + \underbrace{\frac{1}{2} \mathbf{B} \cdot \hat{\mathbf{L}}_G}_{\text{Orbital-Zeeman}} + \underbrace{\frac{1}{8} (\mathbf{B} \times \mathbf{r}_G)^2}_{\text{Diamagnetic}} \quad (1)$$

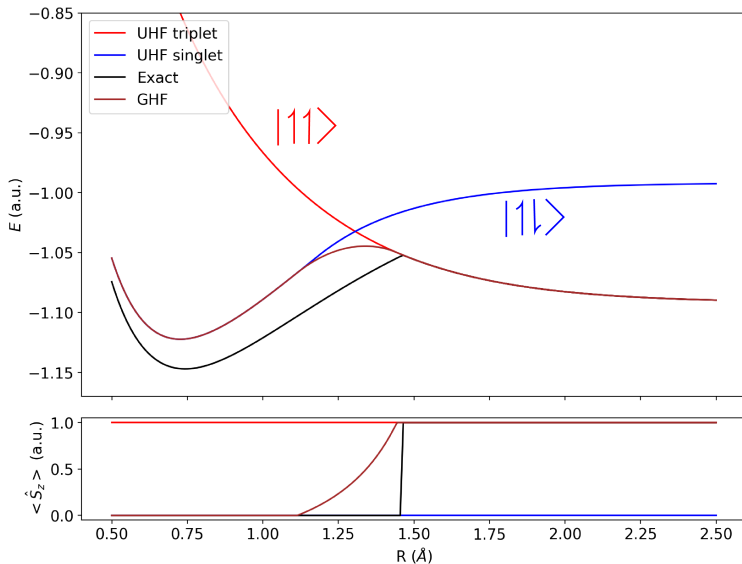
ESTIMATING THE SPIN-FLIP WITH DIFFERENT METHODS



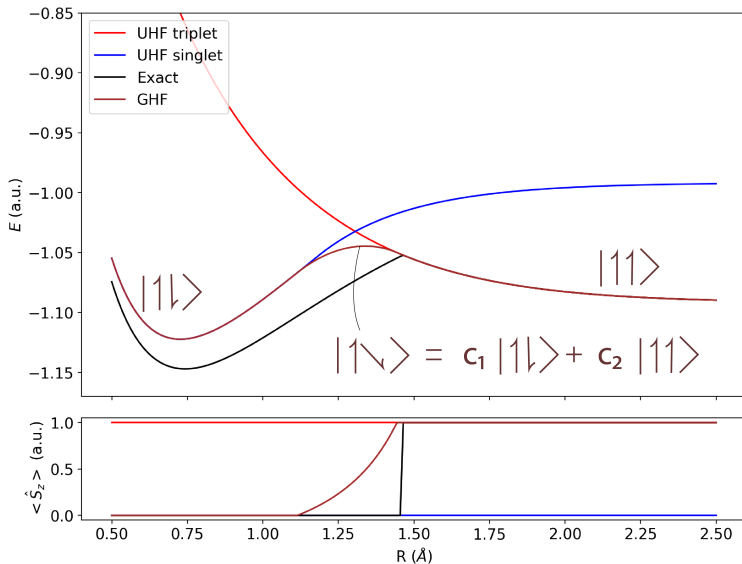
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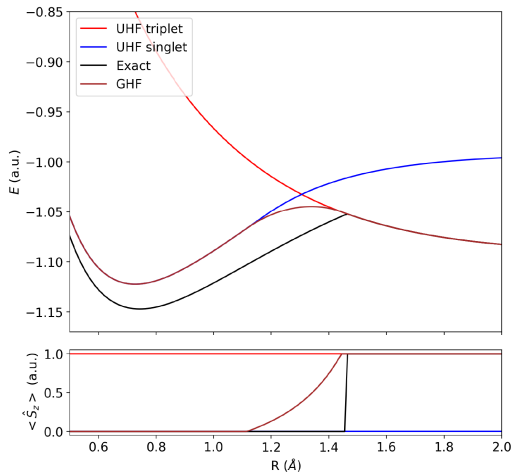
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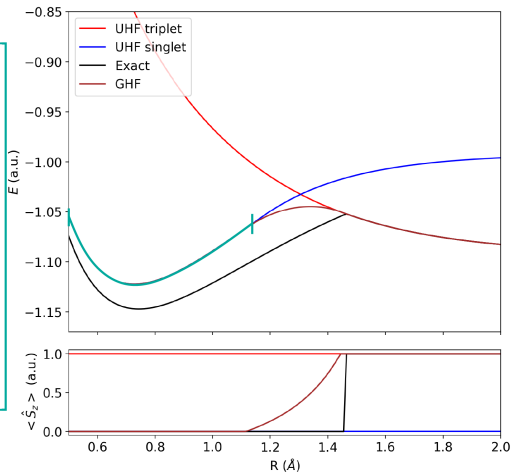
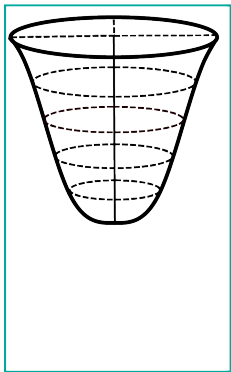
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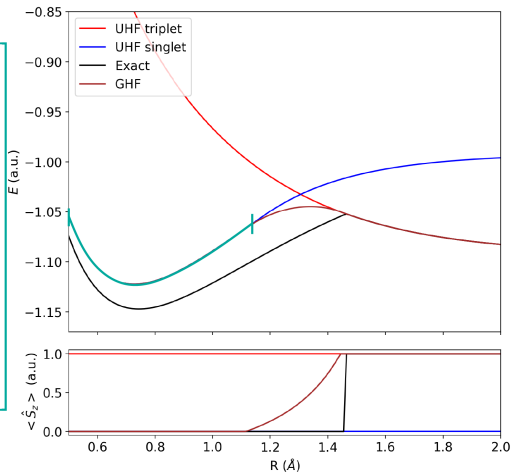
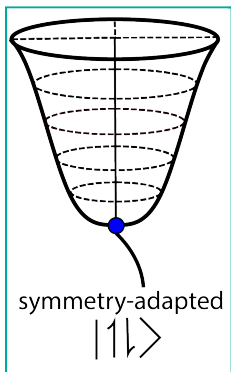
UNDERSTANDING SPONTANEOUS SYMMETRY BREAKING: THE SYMMETRY-BREAKING POTENTIAL



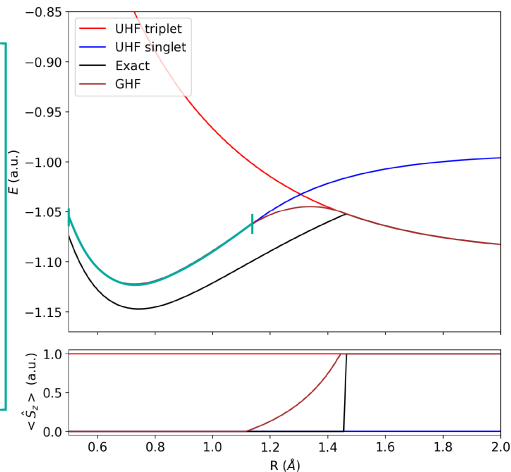
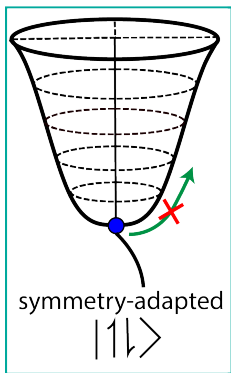
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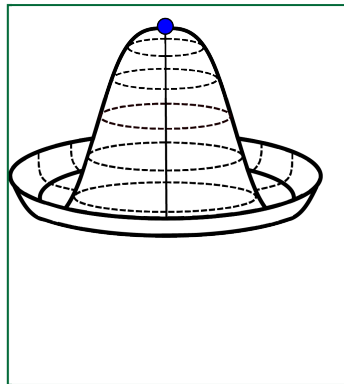
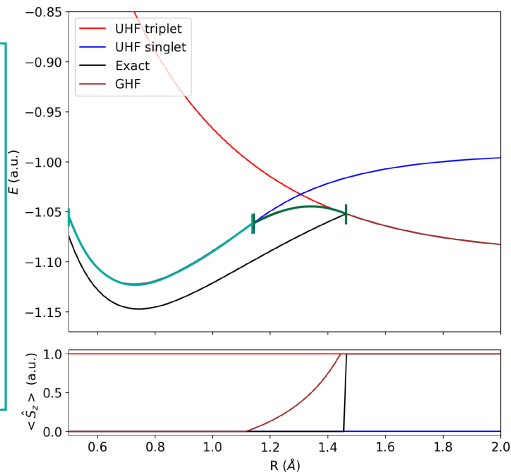
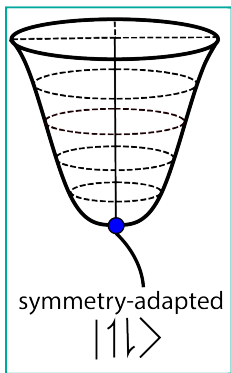
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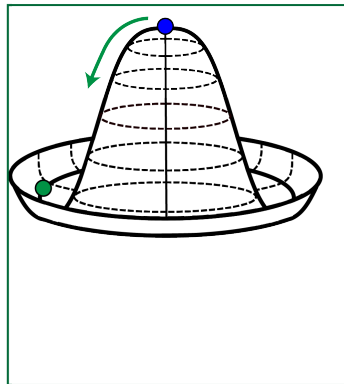
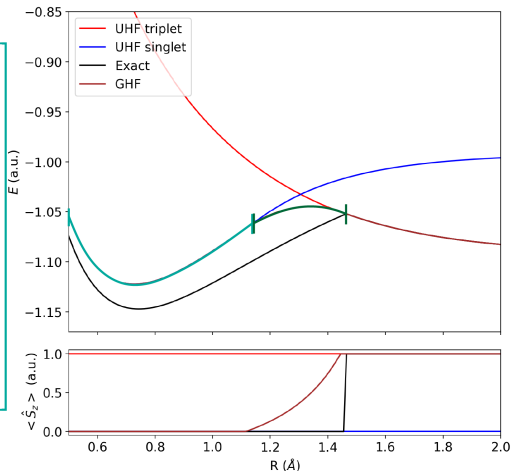
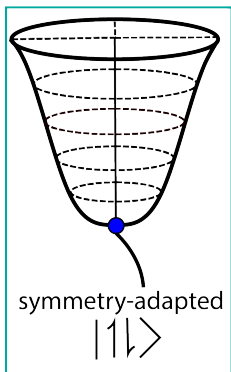
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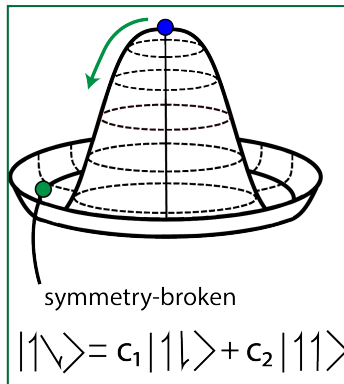
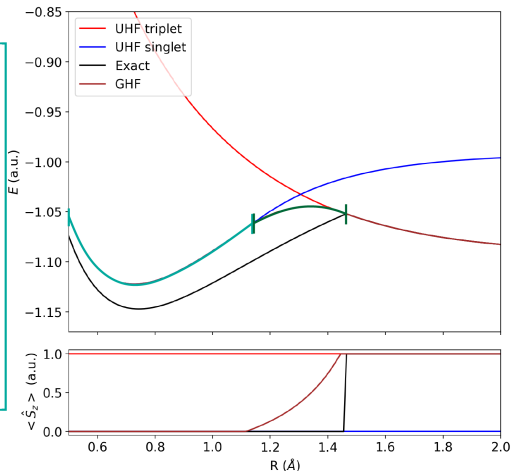
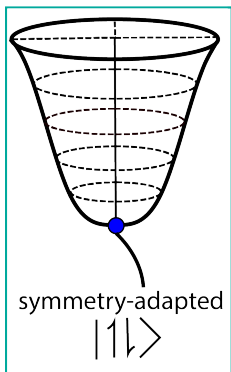
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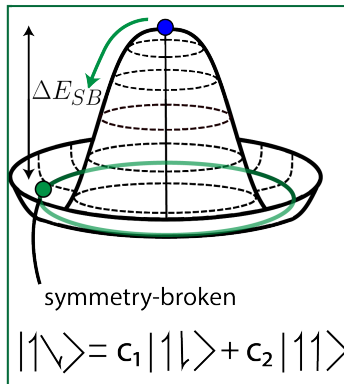
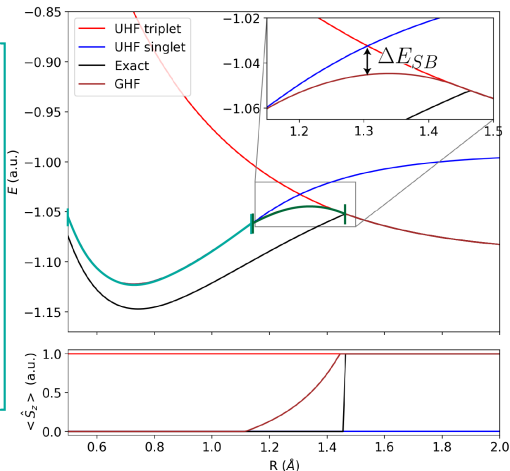
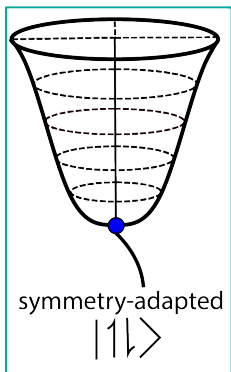
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- The molecular **field-dependent Hamiltonian** (\mathbf{B} along z)

$$\hat{\mathcal{H}}(\mathbf{B}) = \hat{\mathcal{H}}_0 + B_z \hat{S}_z + \frac{1}{2} B_z \hat{L}_{z,G} + \frac{1}{8} B_z^2 (\hat{x}_G^2 + \hat{y}_G^2) \quad (2)$$

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- We obtain the **GHF** state $|\Phi\rangle$ by solving the Roothan-Hall equations in spinor London Atomic Orbital (LAO) basis

$$\mathbf{F} \mathbf{C} = \mathbf{S} \mathbf{C} \epsilon. \quad (4)$$

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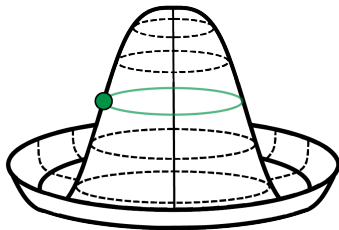
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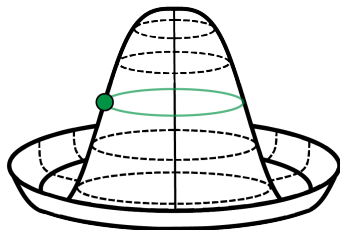
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ACCESSING THE GHF SYMMETRY-BREAKING POTENTIAL

- We introduce the \hat{S}_z -level-shifted Hamiltonian

$$\hat{\mathcal{H}}_{\text{LS}}^{S_z}(\mathbf{B}, \mu) = \hat{\mathcal{H}}(\mathbf{B}) - \mu \hat{S}_z, \quad (5)$$



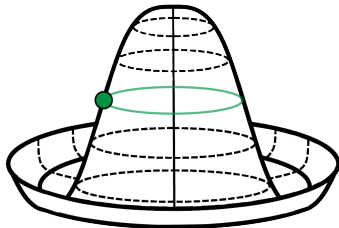
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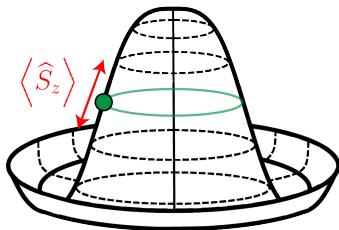
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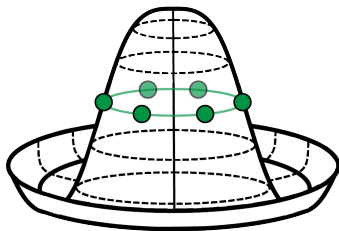
- Now we can **tune** $\langle \hat{S}_z \rangle$ of the **LS-GHF** state $|\Phi_{\text{LS}}^{S_z}(\mu)\rangle$.



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- From a given reference state $|\Phi\rangle$, we can **access** the corresponding **Goldstone manifold** through **spin-rotations**

$$\hat{U}_{R,z}(\theta) |\Phi\rangle = \exp\left(i\theta\hat{S}_z\right) |\Phi\rangle . \quad (7)$$



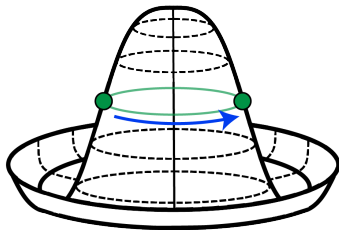
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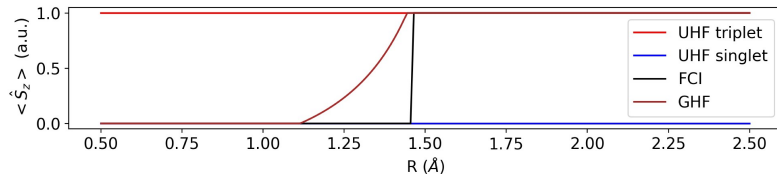
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- We rotate to the **other side** of the Goldstone manifold

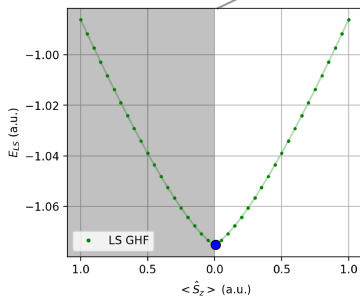
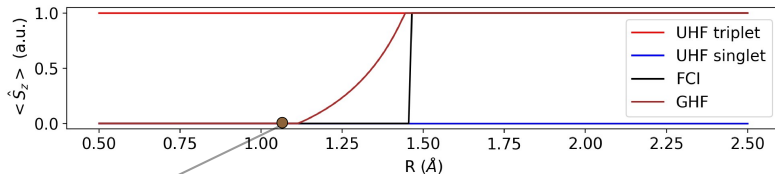
$$\hat{U}_{R,z}(\theta = \pi) |\Phi\rangle = -i\hat{\sigma}_z |\Phi\rangle . \quad (8)$$



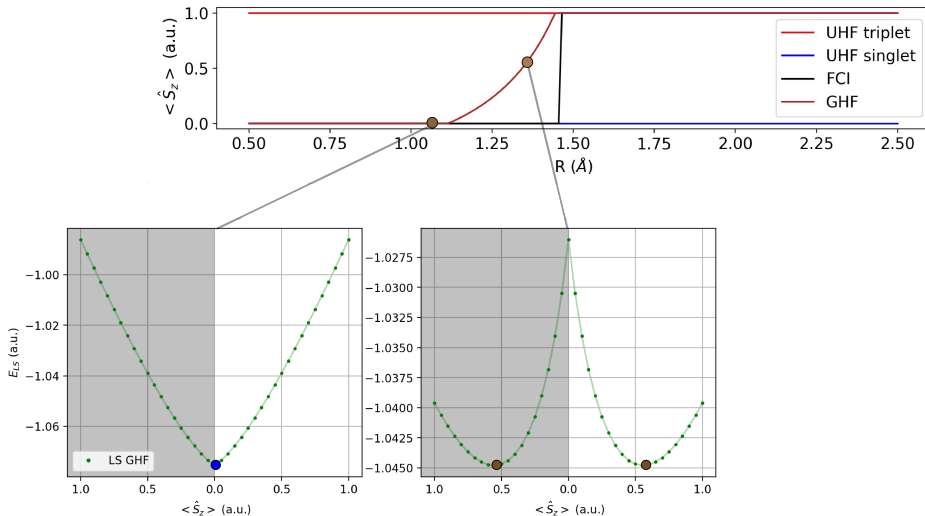
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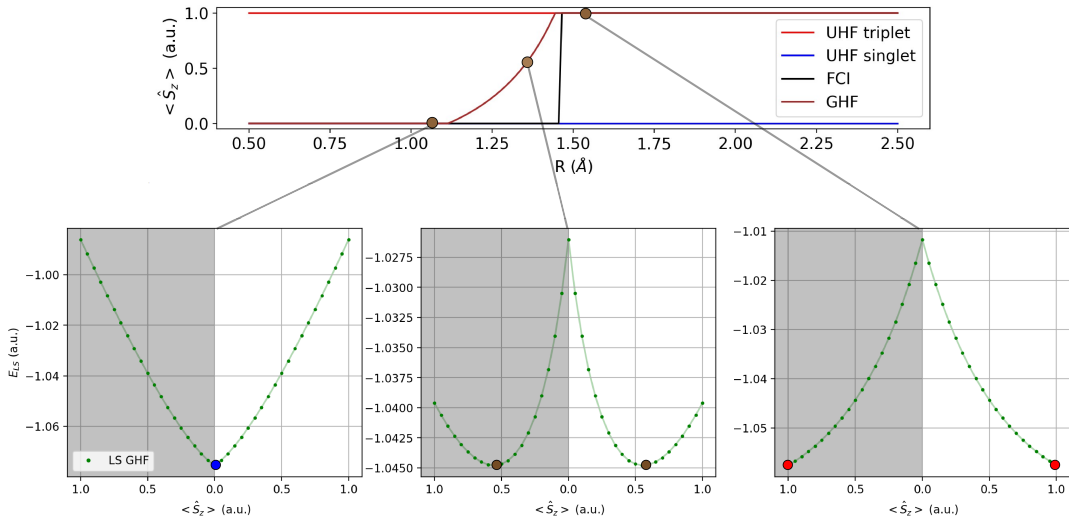
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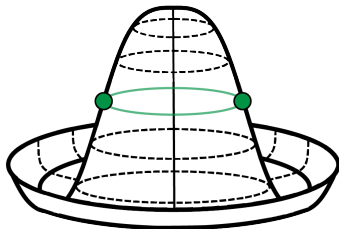


PARTIAL SYMMETRY RESTORATION BY MIXING SYMMETRY-BROKEN STATES

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- We generate a minimal **symmetry-adapted basis** from a level-shifted reference determinant,

$$|\Phi_{\text{LS}}^{S_z}(\mu)\rangle \longrightarrow \{ |\Phi_{\text{LS}}^{S_z}(\mu)\rangle, -i\hat{\sigma}_z |\Phi_{\text{LS}}^{S_z}(\mu)\rangle \}. \quad (9)$$



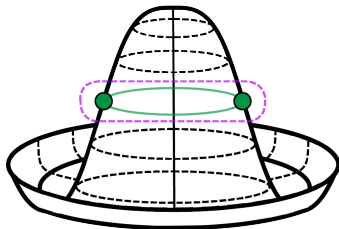
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- Diagonalising $\hat{\mathcal{H}}(\mathbf{B})$ in this basis gives a **partial symmetry-restored ansatz**

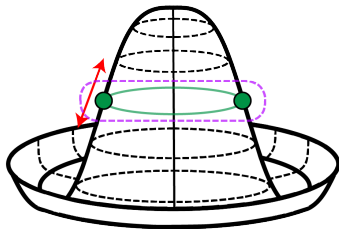
$$|\Psi_i^{S_z}(\mu)\rangle = c_{i,1} |\Phi_{\text{LS}}^{S_z}(\mu)\rangle + c_{i,2} (-i\hat{\sigma}_z |\Phi_{\text{LS}}^{S_z}(\mu)\rangle). \quad (10)$$



VARIATIONALLY TUNING THE UNDERLYING SYMMETRY-BREAKING: LS-VAP

- The LS-VAP ansatz energy is μ -dependent,

$$E_{\Psi}(\mathbf{B}, \mu) = \langle \Psi_0^{S_z}(\mu) | \hat{\mathcal{H}}(\mathbf{B}) | \Psi_0^{S_z}(\mu) \rangle. \quad (11)$$



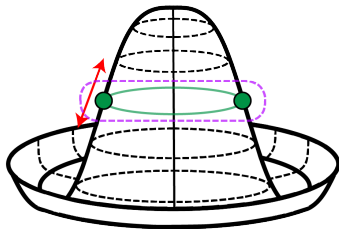
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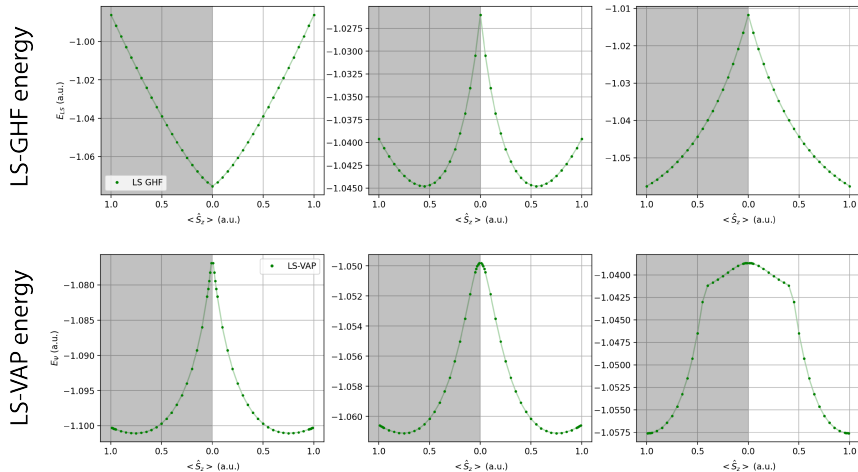
$$E_{\Psi}(\mathbf{B}, \mu) = \langle \Psi_0^{S_z}(\mu) | \hat{\mathcal{H}}(\mathbf{B}) | \Psi_0^{S_z}(\mu) \rangle. \quad (11)$$

- We can tune the height of the sampled Goldstone manifold through variationally optimising the ansatz energy,

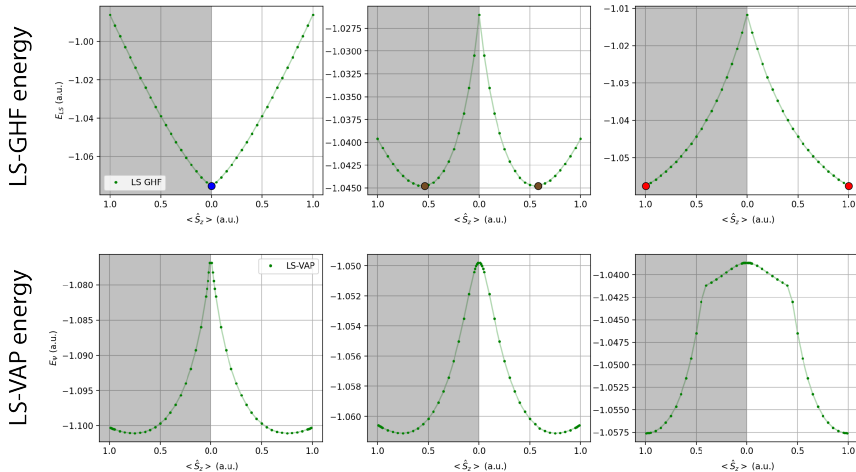
$$\left. \frac{\partial E_{\Psi}(\mathbf{B}, \mu)}{\partial \mu} \right|_{\mu^*} = 0. \quad (12)$$



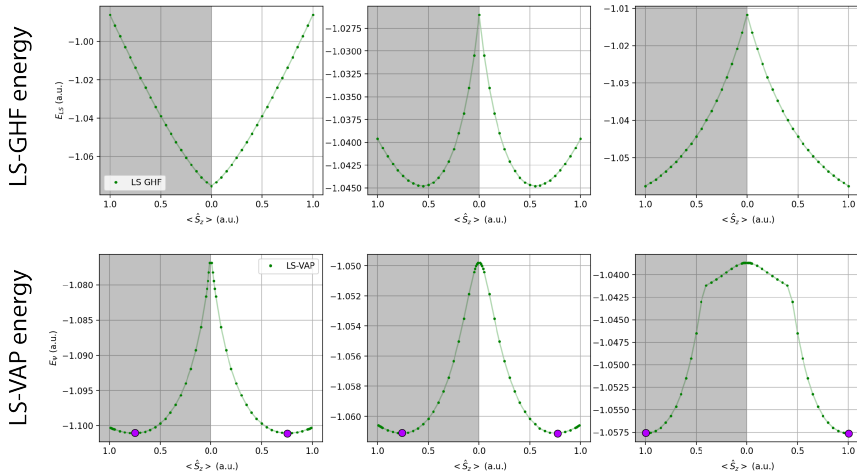
VARIATIONALLY TUNING THE UNDERLYING SYMMETRY-BREAKING: LS-VAP



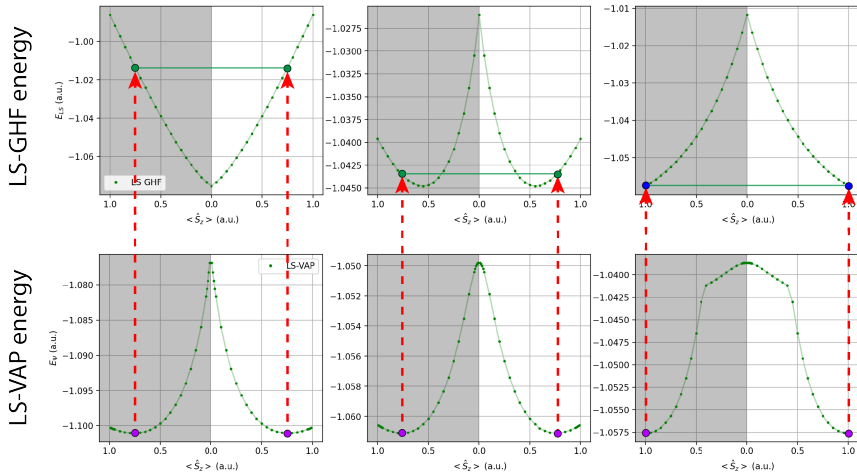
VARIATIONALLY TUNING THE UNDERLYING SYMMETRY-BREAKING: LS-VAP



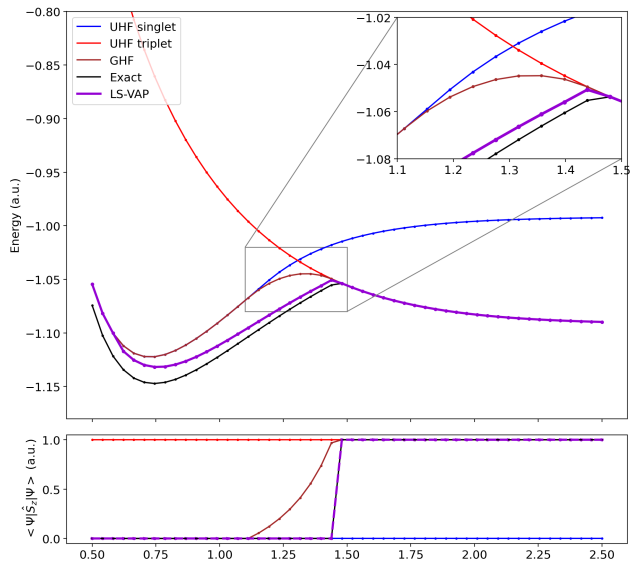
VARIATIONALLY TUNING THE UNDERLYING SYMMETRY-BREAKING: LS-VAP



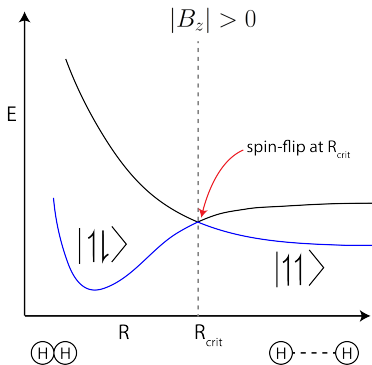
VARIATIONALLY TUNING THE UNDERLYING SYMMETRY-BREAKING: LS-VAP



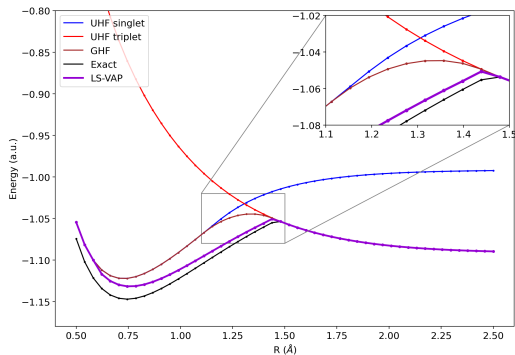
LS-VAP PROVIDES AN ACCURATE SPIN-TRANSITION PREDICTION



CONCLUSIONS



Spin-transitions cause unexpected
**discontinuities throughout
dissociation**



Symmetry-breaking can be **exploited** to gain
dynamic correlation and **accurately describe
spin-transitions**